

### An implemention of the X-FEM for eulerian solid-mechanics

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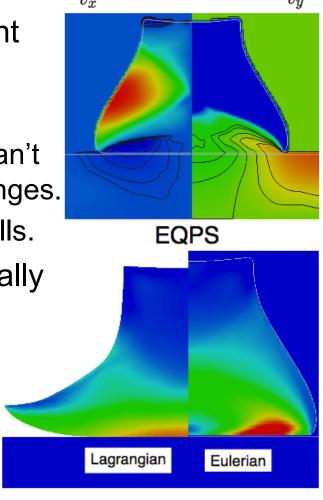
T. Voth, J. Mosso and J. Niederhaus tevoth@sandia.gov

Sandia National Laboratories, NM



### Previous approaches for mixed material cells are problematic for some physics:

- Multi-material problems with significant vorticity/distortion:
  - Lagrangian approaches tangle.
  - Arbitrary Lagrangian Eulerian (ALE) can't merge materials without topology changes.
  - Eulerian approach produces mixed-cells.
- Current mixed-cell approaches generally assume materials are "well" mixed:
  - Assume "equilibrated" state
  - Single velocity/displacement field
  - Lack of intra-element interfaces



### Lagrangian step requires closure model(s) for mixed-cell properties:

$$\mathbf{a}^{n} = \mathbf{M}^{-1} \begin{bmatrix} \mathbf{f}_{hg}^{n} + \mathbf{f}_{ext}^{n} - \int_{\Omega} \mathbf{B}^{t} \boldsymbol{\sigma}^{n} \end{bmatrix}$$

$$\mathbf{v}^{n+1/2} = \mathbf{v}^{n-1/2} + \overline{\Delta t} \mathbf{a}^{n} \quad \mathbf{x}^{n+1} = \mathbf{x}^{n} + \overline{\Delta t} \mathbf{v}^{n+1/2}$$

$$\mathbf{D}^{n+1/2} = \frac{1}{2} (\mathbf{L}^{t} + \mathbf{L})^{n+1/2} \quad \mathbf{D}_{m} = \mathcal{F} (\mathbf{D}, etc...)$$

$$\boldsymbol{\sigma}_{m}^{n+1} = \mathcal{M}_{m} (\boldsymbol{\sigma}_{m}^{n}, \mathbf{D}_{m}^{n+1/2}, etc...) \quad e^{n+1} = e^{n} + \overline{\Delta t} m^{-1} \int_{\Omega} \boldsymbol{\sigma}_{m}^{n} : \mathbf{D}_{m}^{n+1/2}$$

$$\boldsymbol{\sigma} = \mathcal{G} (\boldsymbol{\sigma}_{m}, etc...)$$

Problem	Expected	Predicted
$\overline{v}$	$\overline{v}$	v



#### Why X-FEM:

- Mechanism for intra-element material interfaces.
- Retains base FEM convergence properties.
- Large literature base for X-FEM in context of large deformation, explicit lagrangian mechanics.
- Beginning to be adapted to "operator-split" multimaterial eulerian solid-mechanics [VB06; DLZRM10]:
  - explicit (central difference) lagrangian solve,
  - followed by data transfer "remap" to "better" mesh.

### Goal is to develop capability for (unmixed) intra-element material physics...

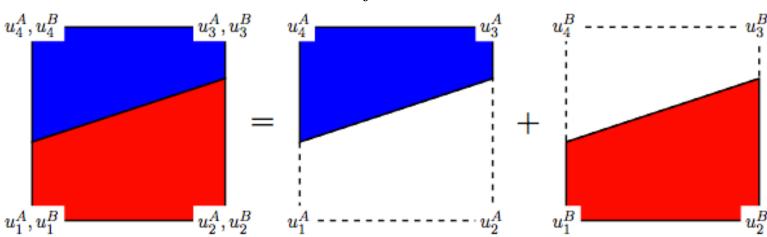
- as a surface phenomenon
- with merging interfaces and fixed mesh topologies
- distinct velocity/displacement fields per material
- while maintaining advantages of current explicitdynamics code-base
- and capitalizing on existing infrastructure.

### We follow the XFEM decomposition approach [HH04, SAB06] ...

Multi-material enriched element equivalent to multiple single-material elements:

$$u^{h}(\mathbf{x}) = \sum_{J} u_{I}^{0}(\mathbf{x}) N_{I}(\mathbf{x}) + \sum_{m} \sum_{J} u_{J}^{m} N_{J}(\mathbf{x}) H_{m}(\mathbf{x})$$
$$u^{h}(\mathbf{x}) = u_{A}^{h}(\mathbf{x}) + u_{B}^{h}(\mathbf{x})$$

$$u^h(\mathbf{x}) = \sum_J u_J^m N_J(\mathbf{x}) H_m(\mathbf{x})$$



### Data transfer can be accomplished in a number of ways:

- With the goal of preserving some key features:
  - Conserve mass, momentum and internal energy.
  - Do not create new minima and maxima (TVD).
  - Volume fractions sum to one after remap.
- Options include:
  - Interpolation (violates conservation) [DLZRM10]
  - Projection methods (violates conservation and TVD)
  - Geometric intersection with Van Leer limiting
    - conservation is built in.
    - limiting controls production of minima and maxima.

# Geometric intersection with Van Leer limiting in two dimensions (1):

Taylor Series provides functional form on donor mesh:

$$f(\mathbf{x}) = \overline{f_e} + (\mathbf{x} - \mathbf{x}_e)^t \mathbf{G}_e + \frac{1}{2} (\mathbf{x} - \mathbf{x}_e)^t \mathbf{H}_e (\mathbf{x} - \mathbf{x}_e)^t - \chi$$
first-
second-
third-order

 $\chi$  provides conservation:

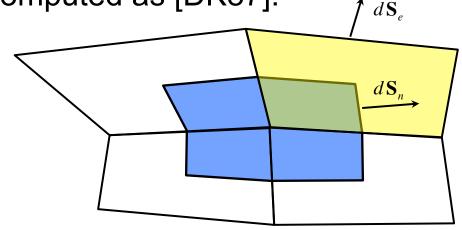
$$\overline{f}_e A_e = \int f(\mathbf{x}) d\Omega_e \Rightarrow \chi = \frac{1}{2A_{e^{-n}}} \int (\mathbf{x} - \mathbf{x}_e)^t \mathbf{H}_e (\mathbf{x} - \mathbf{x}_e^t) d\Omega_e$$

• Gradients/hessians computed as [DK87]:

$$\mathbf{G}_n = \frac{1}{A_n} \oint \overline{f}_e \ d\mathbf{S}_n$$

$$\mathbf{G}_e = \frac{1}{A_e} \sum_{n} A_{n \cap e} \mathbf{G}_n$$

$$\mathbf{H}_e = \frac{1}{A_e} \oint \mathbf{G}_n \, d\mathbf{S}_e$$



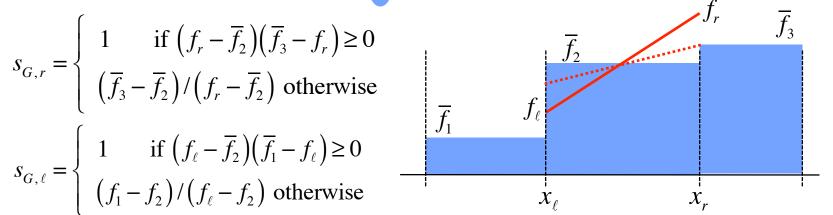
### Geometric intersection with Van Leer limiting in two dimensions (2):

Scale gradient to enforce monotonicity:

$$f(\mathbf{x}) = \overline{f}_e + \mathbf{s}_G (\mathbf{x} - \mathbf{x}_e)^t \mathbf{G}_e$$

$$s_{G,r} = \begin{cases} 1 & \text{if } (f_r - \overline{f}_2)(\overline{f}_3 - f_r) \ge 0\\ (\overline{f}_3 - \overline{f}_2)/(f_r - \overline{f}_2) & \text{otherwise} \end{cases}$$

$$s_{G,\ell} = \begin{cases} 1 & \text{if } (f_{\ell} - \overline{f_2})(\overline{f_1} - f_{\ell}) \ge 0\\ (f_1 - f_2)/(f_{\ell} - f_2) & \text{otherwise} \end{cases}$$



$$s_G = \max\left(\min(s_{G,\ell}, s_{G,r}), 0\right)$$

If third-order also scale hessian terms:

$$f(\mathbf{x}) = \overline{f}_e + s_G (\mathbf{x} - \mathbf{x}_e)^t \mathbf{G}_e + \frac{s_H}{s_H} \left[ (\mathbf{x} - \mathbf{x}_e)^t \mathbf{H}_e (\mathbf{x} - \mathbf{x}_e)^t - \chi \right]$$

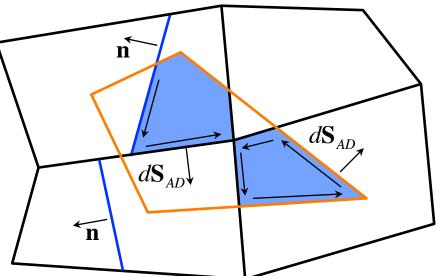
# Geometric intersection with Van Leer limiting in two dimensions (3):

 Integrate function over donor-acceptor intersection element and accumulate to acceptor [D83].

$$\overline{f}_{A} = \frac{1}{A_{A}} \sum_{D} \int_{A_{D} \cap A_{A}} f_{D}(\mathbf{x}) dA = \frac{1}{A_{A}} \sum_{D} \oint_{\Gamma_{AD}} \mathbf{g}_{D}(\mathbf{x}) d\mathbf{S}_{AD}$$

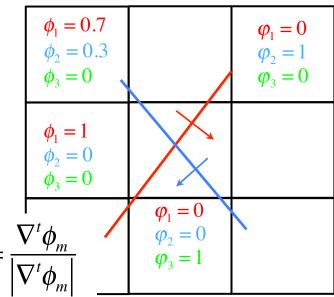
$$\nabla^{t} \mathbf{g}_{D}(\mathbf{x}) \equiv f_{D}(\mathbf{x})$$

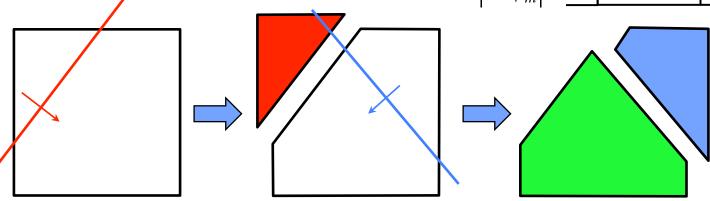
 Further restrict integral to filled region of donor mesh.



### Use interface reconstruction rather than level-set approach: [DVMR08]

- Interfaces rebuilt after remap step.
- Using VOF approach:
  - Compute material volume-fraction gradients.
- Reposition interface along normal to match volume.
  - Remove material from cell.





# Explicit Central Difference discretization requires care for stability ...

$$\mathbf{Ma}^n = \mathbf{f}_{int}^n + \mathbf{f}_c$$

 lumped mass matrix with uniform partitioning of element mass to nodes [MRMCB08]

$$\mathbf{M}^e = \left(\rho^e A^e V_f^e / 4\right) \mathbf{I}_{8\times 8}$$

• matched with gradient operator mean quadrature [SAB06]  $\mathbf{\bar{B}} = \int_{\Omega_e} \mathbf{B} \ d\Omega / A_e$ 

and constraint enforcement between XFEM interfaces:

$$\mathbf{f}_{c} = ?$$

### We understand the issues but use "nodesegment-like" lagrange multipliers...

- ... in an attempt to:
  - minimize interpenetration of XFEM interfaces,
  - and retain a finite stable time-step.
- Other options for explicit XFEM:
  - Merge (small time step) [VB06]
  - Mortar lagrange multiplier (not LBB for XFEM)
  - Penalty (overlap, mass modifications) [DLZRM10]
  - Nitsche's (overlap, mass modifications) [AHD11]
  - Vital Vertices LM (quad robustness) [BMW09, HAD11]
- ... so we use it anyway for it's practicality and economy.

### Forward Increment Lagrange Multiplier approach [CTK91] ...

Algorithm:
$$\mathbf{v}_{0}^{n+1/2} = \mathbf{M}^{-1}\mathbf{f}^{n}$$

$$\lambda_{i+1} = \lambda_{i} + \mathbf{H}\mathbf{r}_{i} \quad \mathbf{H} \approx (\mathbf{G}\mathbf{M}^{-1}\mathbf{G}^{t})$$

$$\mathbf{v}_{i+1}^{n+1/2} = \mathbf{M}^{-1}(\mathbf{f}^{n} - \Delta t\mathbf{G}^{t}\lambda_{i+1})$$

Algorithm:

$$\mathbf{v}_0^{n+1/2} = \mathbf{M}^{-1}\mathbf{f}^n$$

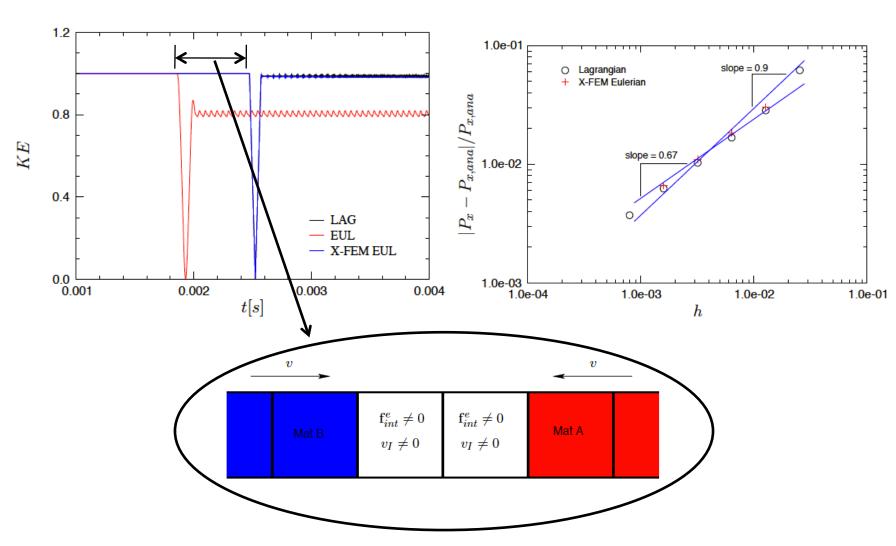
$$\mathbf{r}_i = \mathbf{G} \mathbf{v}_i^{n+1/2}$$

$$\lambda_{i+1} = \lambda_i + \mathbf{H} \mathbf{r}_i \quad \mathbf{H} \approx (\mathbf{G} \mathbf{M}^{-1} \mathbf{G}^t)$$

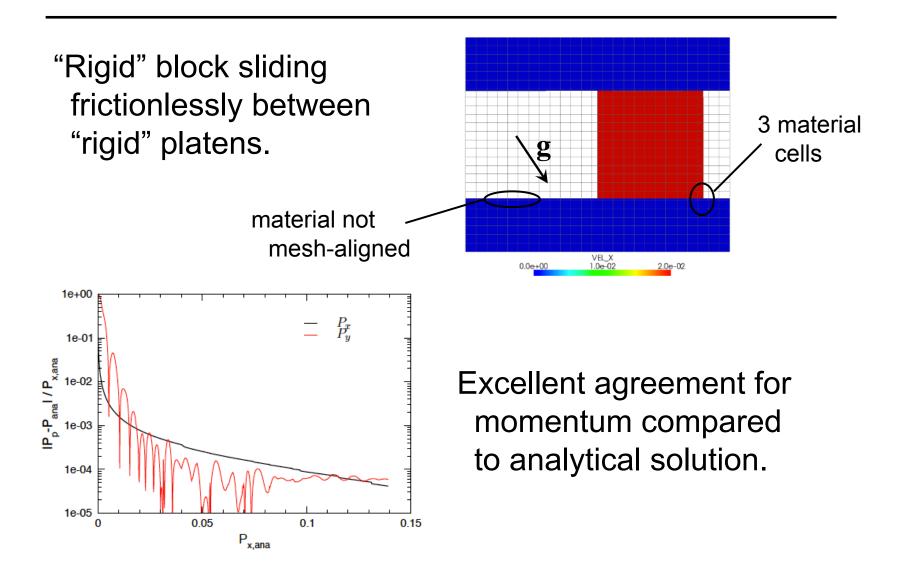
$$\mathbf{v}_{i+1}^{n+1/2} = \mathbf{M}^{-1} \left( \mathbf{f}^n - \Delta t \mathbf{G}^t \lambda_{i+1} \right)$$

No additional limitations to stable time step [DLZRM10,VMR10]

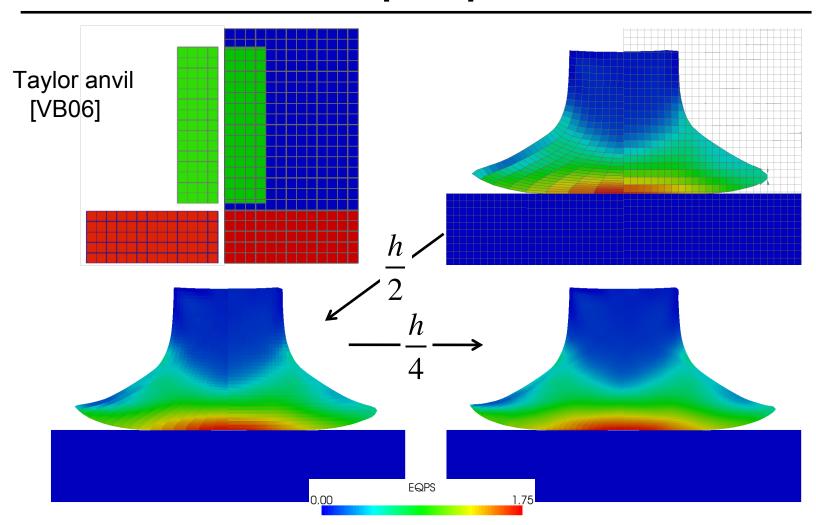
### Provides enhanced results for simple onedimensional problem [CTK91]...



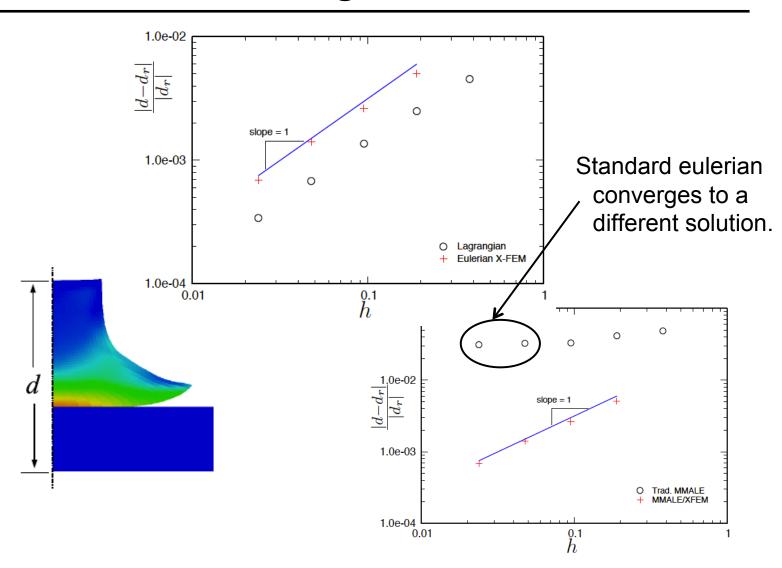
#### ... and simple two-dimensional problems ...



### and accuracy comparable to Lagrangian for more complex problems ...

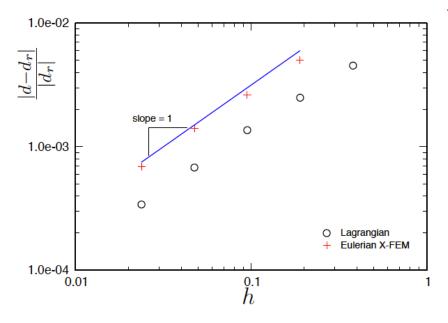


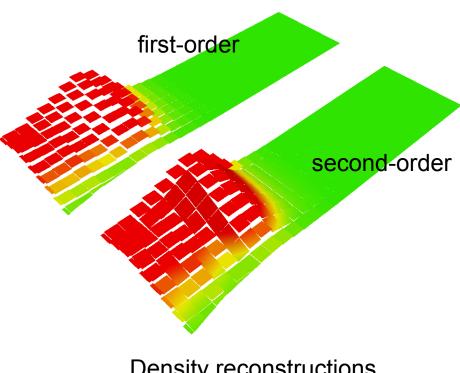
### ... as well as comparable rates of convergence ...



# For this problem remap order has little effect on accuracy/rate-of-convergence ...

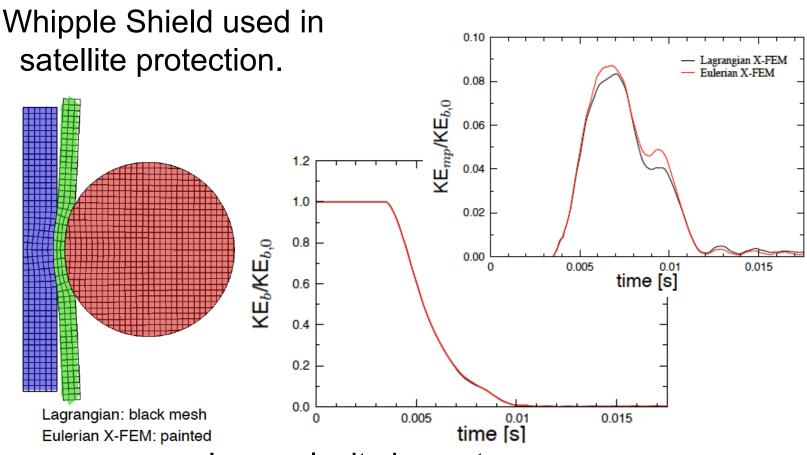
Convergence in height for firstand second-order remap.





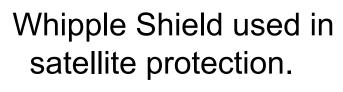
Density reconstructions at early time.

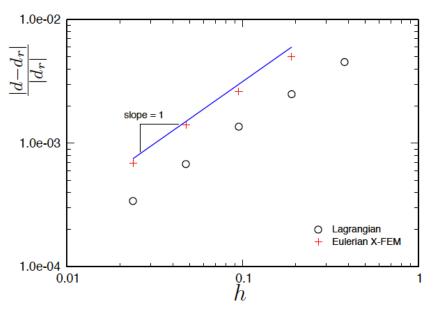
### More complicated problems demonstrate the utility/advantages of the approach ...

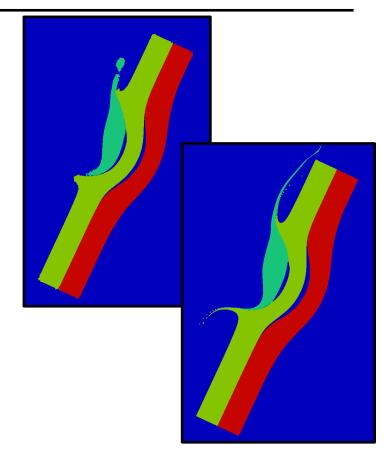


Low-velocity impact preditions compare well.

# More complicated problems demonstrate the utility/advantages of the approach ...







High-velocity impact difficult for lagrangian and unrealistic for eulerian are possible with XFEM.



#### **Conclusions:**

- Developing capability to more accurately treat multimaterial cells in an "operator-split" ALE context.
- Capability builds on existing ALE infrastructure.
- Uses XFEM ideas to provide provide unique kinematics for each material in a cell.
- Uses interface reconstruction rather than level-set ideas to address conservation and complex interface intersections.
- Employs higher-order, conservative remapping algorithms. Advantages are unclear at this point.
- Demonstrates good convergence/accuracy for problems investigated here.



#### (Incomplete) References:

AHD11: Annavarapu et al., IJNME, submitted.

BMW09: Bechet et al., IJNME 78, 931.

CKT91: Carpenter et al., IJNME 32, 103.

D83: Dukowicz, JCP 54, 411.

DK87: Dukowicz and Kodis, SIAM J. Sci Stat Comput 8, 305.

DLZRM10: Dubois et al., Comp Mech 46, 329.

DMRV08: Dolbow et al., CMAME 197, 439.

HAD11: Hautefeuille et al., IJNME, in revision.

HH04: Hansbro and Hansbro, CMAME 193, 3523.

MRMCB08: Menouillard et al., IJNME 74, 447.

SAB06: Song et al., IJNME 67, 868.

VB06: Vitali and Benson, IJNME 67, 1420.

VMR10: Voth et al., USNCCM10, Columbus.